

AS2001: Stellar Structure and Evolution

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Room 316

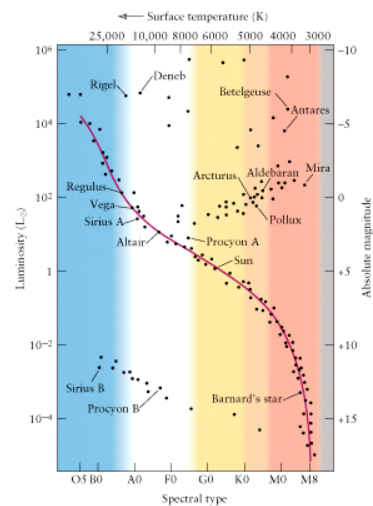
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THE HR DIAGRAM

- Notice **specific** regions:
 - main sequence
 - red giants
 - white dwarfs
- Also - **links** between regions
- and several **empty** regions
- **luminosity range**
 - $10^5 L_{\text{Sun}} - 10^{-4} L_{\text{Sun}}$
- **surface temp range**
 - $10^5 \text{ K} - 2 \times 10^3 \text{ K}$
- **radius range**
 - $10^3 R_{\text{Sun}} - 10^{-2} R_{\text{Sun}}$



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Quick reminder

- The luminosity of a star is $L = 4 \pi R^2 \sigma T^4$
- hence dividing by the luminosity of the Sun

$$\frac{L}{L_{\text{Sun}}} = \left(\frac{R}{R_{\text{Sun}}} \right)^2 \left(\frac{T}{T_{\text{Sun}}} \right)^4$$

$$\text{ie } \log\left(\frac{L}{L_{\text{Sun}}}\right) = 4\log\left(\frac{T}{T_{\text{Sun}}}\right) + 2\log\left(\frac{R}{R_{\text{Sun}}}\right)$$

So, on the HR diagram which shows **log(L)** versus **log(T)**, we can draw a series of straight lines for different stellar radii.

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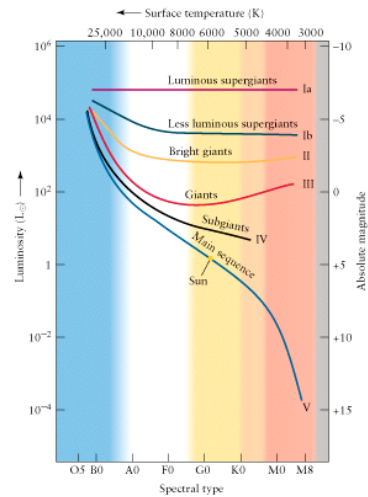
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- For stars of same spectral type (ie surface temp) , the width of the spectral lines varies with the **luminosity** (eg H lines narrow for supergiants, broader for m-s stars).
- This effect is due to variations in the pressure and density of the stellar atmospheres.
- **High pressure** and density \rightarrow many collisions \rightarrow shift in energy levels of H atoms \rightarrow **broad lines**

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- Hence **narrow lines** in supergiants, since their pressure and density is **low** (their mass is spread over a large volume) → few collisions → energy levels undisturbed.
- Hence **luminosity classes I - V**
- **Spectral type + luminosity class** tells you a great deal about a star!



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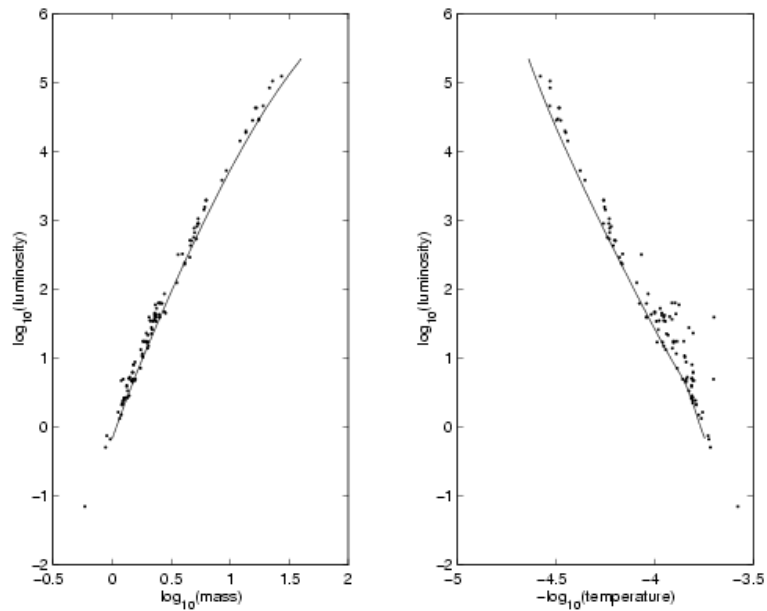
Mass - luminosity relation

- Use **binary systems** to measure stellar masses
- from brightness variations (light curves) of eclipsing binaries, determine orbital inclination i , and sizes of stars relative to their orbit size
- from Doppler effect on positions of spectral lines, measure orbital motion of both stars, and hence absolute size of orbits (e.g. in km).
- combined data give masses, radii, temperatures and luminosities of both stars, usually quoted in solar units.

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Mass-luminosity and HR diagram (T-L) for detached eclipsing binaries



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- From masses of many stars, find for m-s stars

$$\frac{L}{L_{\text{Sun}}} = \left(\frac{M}{M_{\text{Sun}}} \right)^{4.0 \pm 0.02} \quad \text{for } 0.4 < M < 10 M_{\text{Sun}}$$

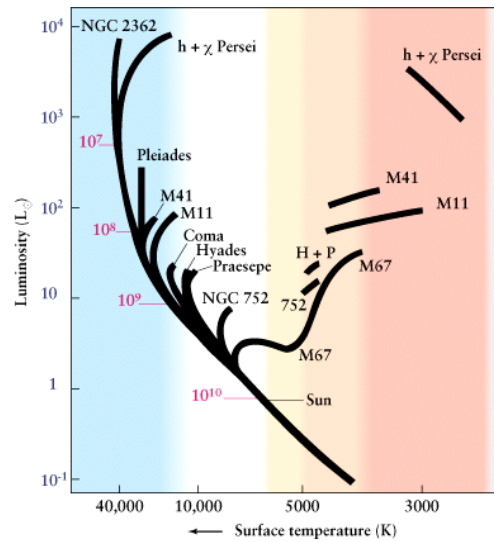
$$\frac{L}{L_{\text{Sun}}} = \left(\frac{M}{M_{\text{Sun}}} \right)^{3.6 \pm 0.1} \quad \text{for } 5 \leq M \leq 40 M_{\text{Sun}}$$

But note exponent in this power law decreases rapidly for stars of higher mass

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- **Composite HR diagrams of different clusters → stellar evolution; also use clusters to get distances**



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Open clusters:

- **distances: main sequence (m-s) fitting**
- **distribution: Milky Way (MW) plane**
- **population: I**
- **Chemical comp. :**
 $X \sim 0.73$; $Y \sim 0.25$; $Z \sim 0.01 - 0.04$
- **Ages: young - old; $10^6 - n \times 10^9$ yrs**

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Globular clusters:

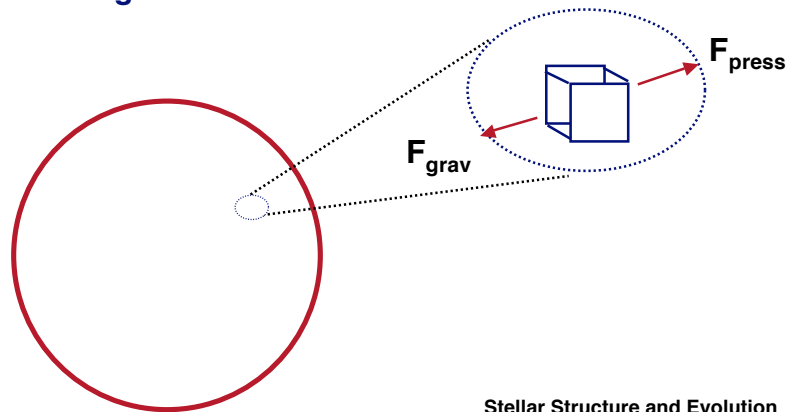
- **distances:** m-s fitting; RR Lyrae (all have $M_V \sim 0.6$)
- **distribution:** galactic halo + nucleus
- **population:** II
- **Chemical comp. :** $Z \sim 0.0001 - 0.001$
- **Ages:** very old; $n \times 10^{10}$ yrs

STELLAR STRUCTURE

- **Governed by mass, age, initial chemical composition**
- **The equations of stellar structure describe the complete internal structure of a star**
- **this in turn determines the observable parameters of the star: its radius, surface temperature and luminosity**
 - **NB: asteroseismology provides further tests of these equations and our knowledge of stellar structure**

Stellar equilibrium

- Most stars are in a stable state (i.e. there is a balance between all their internal forces)
- the inward force of **gravity** is balanced by an outward **pressure** force (hydrostatic equilibrium)
- pressure is higher in centre of star



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Gravitational force

Newton's Laws:

- the gravitational force F_{grav} is the same as if all the mass closer to the star than the element were concentrated at the centre.
- The outer layers have no effect.

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- If each side of the cube is of area A , it is a distance r from the centre and has density ρ then the mass of the element is $m = \rho A dr$ and the mass interior to r is $M(r)$, giving



$$F_{\text{grav}} = \frac{GM(r)\rho A dr}{r^2}$$

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Pressure force

- Pressure forces all balance **except** on the lower and upper faces
- lower face: pressure force outward $P(r)A$
- upper face: pressure force inward $P(r + dr)A$
- condition for no net force is:

$$P(r + dr)A - P(r)A + \frac{GM(r)\rho A dr}{r^2} = 0$$

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- **But, for an infinitesimal element of side dr ,**

$$P(r + dr) - P(r) = \left(\frac{dP}{dr} \right) dr$$

- **and so force balance becomes**

$$\frac{dP(r)}{dr} = - \frac{GM(r)\rho(r)}{r^2} \quad (1)$$

The equation of hydrostatic equilibrium

- **The mass interior to r can be found from**

$$\frac{dM(r)}{dr} = 4\pi r^2 \rho(r) \quad (2)$$

i.e.

$$M(r) = 4\pi \int r^2 \rho(r) dr$$