## 12. Circumstellar Matter Monte Carlo Radiation Transfer I

- Monte Carlo "Photons" and interactions
- Sampling from probability distributions
- Optical depths, isotropic emission, scattering


## Monte Carlo Basics

- Emit energy packet, hereafter a "photon"
- Photon travels some distance
- Something happens...

- Scattering, absorption, re-emission


## Photon Packets

If the total input luminosity is $L$, then each photon packet carries energy $E_{i}=L \Delta t / N$, where $N$ is the number of Monte Carlo photons. A Monte Carlo photon represents $N_{\gamma}$ real photons, where $N_{\gamma}=E_{i} / h v_{i}$. So, a Monte Carlo photon packet moving along direction specified by $\theta$ will contribute to the specific intensity:


Note, $I_{\mathrm{v}}$ is a distribution function. We will be working with discrete energies. Binning the photon packets into directions, frequencies, etc, enables us to simulate a distribution function. e.g., spectrum: bin in frequency; scattering phase function: bin in angle


## Photon Interactions

Volume $=A \mathrm{~d} l$

Number density $n$
Cross section $\sigma$
A
$\mathrm{d} l$
Energy removed from beam per $t / \mathrm{V} / \mathrm{d} \Omega=I_{v} \sigma$

Number of photons absorbed/scattered from beam per sec

$$
=I_{v} \sigma n A \mathrm{~d} l
$$

Number of photons absorbed/scattered from beam per sec per area

$$
=I_{v} \sigma n \mathrm{~d} l
$$

Intensity differential over $\mathrm{d} l$ is $\mathrm{d} I_{v}=-I_{\mathrm{v}} n \sigma \mathrm{~d} l$. Therefore

$$
I_{v}(l)=I_{v}(0) \exp (-n \sigma l)
$$

Fraction scattered or absorbed $/$ length $=n \sigma$
$n \sigma=$ volume absorption coefficient $=\rho \kappa$
Mean free path $=1 / \mathrm{n} \sigma=$ average distance between interactions
Probability of interaction over $\mathrm{d} l$ is $n \sigma \mathrm{~d} l$
Probability of traveling $\mathrm{d} l$ without interaction is $1-n \sigma \mathrm{~d} l$

$N$ segments of length $L / N$
Probability of traveling $L$ before interacting is

$$
\begin{aligned}
\mathrm{P}(L) & =(1-n \sigma l / N)(1-n \sigma l / N)(1-n \sigma l / N) \ldots \\
& =(1-n \sigma l / N)^{N}=\exp (-n \sigma L) \\
\mathrm{P}(L) & =\exp (-\tau)
\end{aligned}
$$

$\tau=$ number of mean free paths over distance $L$.

## Probability Distribution Function

The probability distribution function (PDF) for photons to travel optical depth $\tau$ before an interaction is $\exp (-\tau)$. If we pick $\tau$ uniformly over the range 0 to infinity we will not reproduce $\exp (-\tau)$. We want to pick lots of small $\tau$ s and fewer large $\tau \mathrm{s}$. Same with a scattering phase function: want to get the correct number of photons scattered into different directions, forward and back scattering, etc.


## Cumulative Distribution Function

$\mathrm{CDF}=$ Area under $\mathrm{PDF}=\int P(x) \mathrm{d} x$

Want to randomly choose $\tau, \theta, \lambda, \ldots$ so that PDF is reproduced
$\xi$ is a random number uniformly chosen in range $[0,1]$

$$
\xi=\int_{0}^{X} P(x) \mathrm{d} x \Rightarrow X
$$

$$
\int_{-\infty}^{\infty} P(x) \mathrm{d} x=1
$$

The above equation is the fundamental principle behind Monte Carlo techniques and is used to sample randomly from PDFs.
e.g., $P(\theta)=\cos \theta$ and we want to map $\xi$ to $\theta$. Choose random $\theta$ s to "fill in" $P(\theta)$


If we sample many random $\theta_{i}$ in this way and "bin" them, we will reproduce the curve $P(\theta)=\cos \theta$.

## Choosing a Random Optical Depth

$P(\tau)=\exp (-\tau)$, i.e., photon travels $\tau$ before interaction

$$
\xi=\int_{0}^{\tau} \mathrm{e}^{-\tau} \mathrm{d} \tau=1-\mathrm{e}^{-\tau} \Rightarrow \tau=-\log (1-\xi)
$$

Since $\xi$ is in range $[0,1]$, then $(1-\xi)$ is also in range $[0,1]$, so we may write:

$$
\tau=-\log \xi
$$

We find the physical distance, $L$, that the photon has traveled from:

$$
\tau=\int_{0}^{L} n \sigma \mathrm{~d} s
$$

## Random Isotropic Direction

Solid angle is $\mathrm{d} \Omega=\sin \theta \mathrm{d} \theta \mathrm{d} \phi$, we want to choose $(\theta, \phi)$ so they fill in PDFs for $\theta$ and $\phi . P(\theta)$ normalized over $[0, \pi], P(\phi)$ normalized over $[0,2 \pi]$ :

$$
P(\theta)=1 / 2 \sin \theta \quad P(\phi)=1 / 2 \pi
$$

Using fundamental principle from above:

$$
\begin{aligned}
& \xi=\int_{0}^{\theta} P(\theta) \mathrm{d} \theta=\frac{1}{2} \int_{0}^{\theta} \sin \theta \mathrm{d} \theta=\frac{1}{2}(\cos \theta-1) \\
& \xi=\int_{0}^{\phi} P(\phi) \mathrm{d} \phi=\frac{1}{2 \pi} \int_{0}^{\phi} \mathrm{d} \phi=\frac{\phi-1}{2 \pi}
\end{aligned}
$$

$$
\begin{aligned}
& \theta=\cos ^{-1}(2 \xi-1) \\
& \phi=2 \pi \xi
\end{aligned}
$$

Use this formula for emitting photons isotropically from a point source, or for choosing a scattering direction for isotropic scattering.

## Rejection Method

The rejection method is used when we cannot invert the PDF as in the above examples to obtain analytic formulae for random $\theta, \lambda$, etc.


Choose $x_{1}$ in range $[a, b]: x_{1}=a+\xi(b-a)$, calculate $P\left(x_{1}\right)$
Choose $y_{1}$ in range $\left[0, P_{\max }\right]: y_{1}=\xi P_{\text {max }}$
If $y_{1}>P\left(x_{1}\right)$, reject $x_{1}$. Choose new $x_{2}, y_{2}$ until $y_{2}<P\left(x_{2}\right)$ : accept $x_{2}$ Efficiency = Area under $P(x)$

## Calculate $\pi$ by the Rejection Method

Choose $N$ random positions $\left(x_{i}, y_{i}\right)$ :


Choose $x_{i}$ in range $[-R, R]: x_{i}=(2 \xi-1) R$ Choose $y_{i}$ in range $[-R, R]: y_{i}=(2 \xi-1) R$ $\operatorname{Reject}\left(x_{i}, y_{i}\right)$ if $x_{i}^{2}+y_{i}^{2}>R^{2}$ Number accepted / $N=\pi R^{2} / 4 R^{2}$

$$
N_{A} / N=\pi / 4
$$

Increase accuracy (signal to noise) by increasing $N$.
FORTRAN 77:

$$
\begin{aligned}
& \text { do } \mathrm{i}=1, \mathrm{~N} \\
& \mathrm{x}=2 .{ }^{*} \operatorname{ran}-1 . \\
& \mathrm{y}=2 .{ }^{*} \operatorname{ran}-1 . \\
& \text { if }\left(\left(x^{*} \mathrm{x}+\mathrm{y}^{*} \mathrm{y}\right) . \text { lt. 1. }\right) \mathrm{NA}=\mathrm{NA}+1 \\
& \text { end do } \\
& \text { pi }=4 .{ }^{*} \mathrm{NA} / \mathrm{N} \\
& \hline
\end{aligned}
$$

## Albedo

When photon gets to interaction location at the randomly chosen optical depth, $\tau$, we must decide whether the photon is scattered or absorbed. We use the albedo or scattering probability. It is the ratio of scattering to total opacity:

$$
a=\frac{\sigma_{S}}{\sigma_{S}+\sigma_{A}}
$$

To decide if a photon is scattered we choose a random number in the range $[0,1]$ and scatter if $\xi<a$, otherwise the photon is absorbed.

We now have the tools required to write a Monte Carlo radiation transfer program for isotropic scattering in a constant density slab...

