

Interconnectedness

Moments (J_ν , H_ν , K_ν) depend on I_ν

Need to solve ERT to get I_ν

I_ν (and hence J_ν) depends on position and direction

I_ν depends on S_ν , hence on emissivity and opacity

Opacity depends on temperature and ionization

Temperature and ionization depends on J_ν

$$J_\nu = \frac{1}{4\pi} \int I_\nu \, d\Omega$$

$$H_\nu = \frac{1}{4\pi} \int I_\nu \cos \theta \, d\Omega$$

$$K_\nu = \frac{1}{4\pi} \int I_\nu \cos^2 \theta \, d\Omega$$

$$\frac{dI_\nu}{d\tau_\nu} = S_\nu - I_\nu$$

$$S_\nu \equiv \frac{j_\nu}{\alpha_\nu}$$

$$d\tau_\nu = \alpha_\nu(s) \, ds = \rho(s) \kappa_\nu \, ds$$

Example: Model H II Region

- Sources of ionizing photons
- Opacity from neutral H: bound-free
- 1st iteration:
 - Medium fully ionized (no neutral H) so opacity is zero
 - Solve ERT throughout medium to get J_ν
 - Solve for ionization structure, some regions neutral
- 2nd iteration:
 - new opacity structure,
 - different solution for ERT, different J_ν values
 - new ionization and opacity structure
- Iterate until get convergence: solution of ERT, J_ν , ionization structure do not change with further iterations

Sources of Opacity

- Bound-bound transitions
- Bound-free transitions
- Free-free transitions (Bremstrahlung)
- Scattering: electrons, molecules, dust,...

Total opacity:

$$\rho \kappa(\nu) = \sum_i n_i \sigma_i(\nu)$$

Bound-Bound Line Transitions

Between upper, u , and lower, l , levels of atom, ion, or molecule may occur as:

Spontaneous Radiative De-excitation

Radiative Excitation

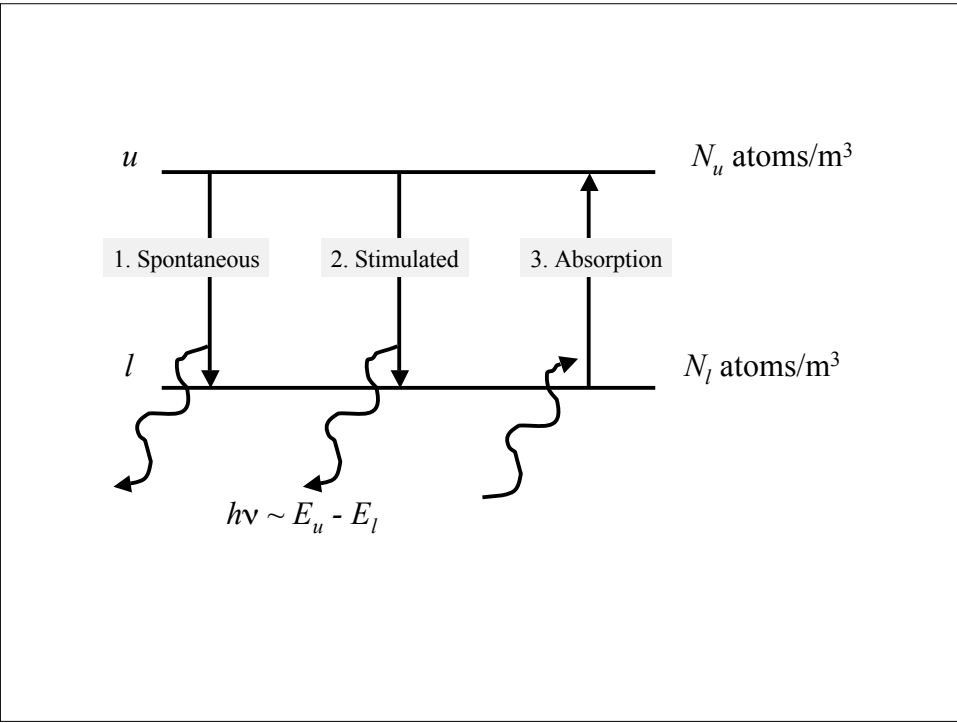
Induced Radiative De-excitation

Collisional Excitation

Collisional De-excitation

Use *Einstein Coefficients* to describe these processes,

$$A_{ul}, B_{ul}, B_{lu}, C_{ul}, C_{lu}$$



Spontaneous De-excitation

$A_{ul} \equiv$ Probability for spontaneous de-excitation from u to l , per second, per particle in state u

Mean lifetime of particles in state u : $\Delta t = 1/A_{ul}$ secs
 Spread in energy (Heisenberg) yields *Lorentz profile*:

$$\psi(\nu - \nu_0) = \frac{\gamma^{\text{rad}} / 4\pi^2}{(\nu - \nu_0)^2 + (\gamma^{\text{rad}} / 4\pi)^2}$$

$\gamma^{\text{rad}} = 1/\Delta t$
 A_{ul} is a summation over the profile

Radiative Excitation

$B_{lu} \bar{J}_{\nu_0}^{\varphi} \equiv$ Number of radiative excitations from l to u , per second, per particle in state l

Induced De-excitation

$B_{ul} \bar{J}_{\nu_0}^{\chi} =$ Number of induced radiative de-excitations from u to l , per second, per particle in state u

Averaged mean intensity over line φ, χ :

$$\bar{J}_{\nu_0}^{\varphi} \equiv \int_0^{\infty} J_{\nu} \varphi(\nu - \nu_0) d\nu; \quad \int_0^{\infty} \varphi(\nu - \nu_0) d\nu = 1$$

Line Profiles

- Line profiles give probabilities for emitting/absorbing photon with frequency ν
- Number of atoms able to absorb photon = $N_l \varphi(\nu - \nu_0)$
- Complete redistribution: no “memory” of preceding processes: incident (exciting) and emitted photons not correlated
- Complete redistribution: profile shapes equal, $\psi = \varphi = \chi$

Collisional Excitation / De-excitation

$C_{lu} \equiv$ Number of collisional excitations from l to u , per second, per particle in state l

$C_{ul} \equiv$ Number of collisional de-excitations from u to l , per second, per particle in state u

Electron collisions usually most important:
Cooling in H II regions: collisions with electrons excite ions => de-excite and emit photons, carry away energy, hence important cooling source

Einstein Relations

Einstein coefficients coupled by the Einstein relations:

$$\frac{B_{lu}}{B_{ul}} = \frac{g_u}{g_l}; \quad \frac{A_{ul}}{B_{ul}} = \frac{2h\nu^3}{c^2}$$
$$\frac{C_{lu}}{C_{ul}} = \frac{g_u}{g_l} e^{E_{lu}/kT}$$

g = statistical weight or degeneracy: $g = 2J + 1$

Line Source Function

$$S_{\nu}^l \equiv j_{\nu}^l / \alpha_{\nu}^l = \frac{n_u A_{ul} \psi(\nu - \nu_0)}{n_l B_{lu} \varphi(\nu - \nu_0) - n_u B_{ul} \chi(\nu - \nu_0)}$$

Einstein relations:

$$S_{\nu}^l = \frac{\frac{A_{ul} \psi}{B_{ul} \varphi}}{\frac{n_l B_{lu} \chi}{n_u B_{ul} \varphi}} = \frac{2h\nu^3}{c^2} \frac{\psi / \varphi}{\frac{g_u n_l \chi}{g_l n_u \varphi}}$$

Complete
Redistribution:

$$S_{\nu_0}^l = \frac{n_u A_{ul}}{n_l B_{lu} - n_u B_{ul}} = \frac{2h\nu_0^3}{c^2} \frac{1}{\frac{g_u n_l}{g_l n_u} - 1}$$

Continuum Transitions

Bound-free:

Extinction cross section for hydrogen and hydrogen-like transitions, **Kramer's** formula:

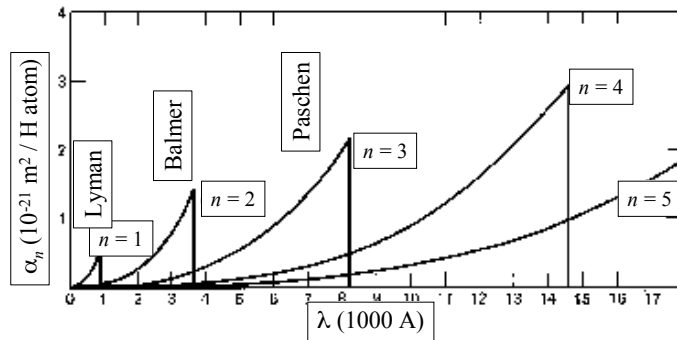
$$\sigma_{\nu}^{\text{bf}} = 2.815 \times 10^{25} \frac{Z^4}{n^5 \nu^3} g_{\text{bf}} \text{ m}^2 \quad \text{for } \nu \geq \nu_0$$

n = principal quantum number of level i from which the atom or ion is ionized.

Z = ion charge and g_{bf} = *Gaunt factor*, QM correction ~ 1

Cross section decays as $1/\nu^3$ above the threshold ("edge") frequency ν_0 and is zero below it.

H I Continuous Extinction



e.g., Paschen continuum: absorptions to $n = 3$

Balmer Jump at $\lambda = 3646 \text{ \AA}$

Peaks increase with n : Rydberg sequence $h\nu_n = 13.6/n^2 \text{ eV}$

Free-free:

Thermal Bremsstrahlung, cross section is:

$$\sigma_v^{\text{ff}} \approx 3.7 \times 10^4 N_e N_{\text{ion}} \frac{Z^2}{T^{1/2} \nu^3} (1 - e^{-h\nu/kT}) g_{\text{ff}} \text{ m}^2$$

Z = ion charge,

N_e, N_{ion} : electron, ion densities

g_{ff} = Gaunt correction factor

Scattering

Electron (Thomson) Scattering:

Frequency independent for low energy photons:

$$\sigma_v^T \equiv \sigma^T = \frac{8\pi}{3} r_e^2 = 6.65 \times 10^{-29} \text{ m}^2$$

Frequency dependent:

High energy photons: *Compton scattering*

High energy electrons: *inverse-Compton scattering*

Thomson scattering is major source of continuous extinction in hot star atmospheres where H is ionized

Rayleigh Scattering:

Cross section for photons with $\nu \ll \nu_0$ by bound electrons with binding energy $h \nu_0$ is:

$$\sigma_v^R \approx f_{lu} \sigma^T \left(\frac{\nu}{\nu_0} \right)^4$$

f_{lu} and ν_0 characterize the major bound-bound “resonance transition” of the bound electron

e.g., Ly α transition in neutral hydrogen of a weighted sum over all Lyman lines

ν^4 ($1/\lambda^4$) dependence makes the sky blue and sunsets red

Redistribution in Angle or Scattering Phase Function:

Thomson and Rayleigh scattering are *coherent*: photon re-directed with same ν

Re-direction has *phase function* $\sim 1 + \cos^2\theta$

This is the angular shape for the scattered photons