

2.4 Energy conservation II: Transport.

- The Sun's interior is hotter than its surface.
- Existence of a temperature gradient implies an outward flux of energy.
- Energy flux is determined by conservation of energy as just shown.
- Temperature gradient depends on method of energy transport:
 - Radiative diffusion
 - Convective motions

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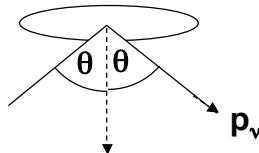
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Flux and radiation pressure

- From Nebulae/Atmospheres, flux through surface element dA in frequency interval from ν to $\nu + d\nu$ is:

$$F_\nu = \oint_{4\pi} I_\nu \cos \theta \cdot d\Omega = 2\pi \int_{-1}^1 I_\nu \mu d\mu$$

where we have substituted $\mu = \cos \theta$.



- Each photon carries momentum $p_\nu = E/c = h\nu/c$
- Bounces off dA at angle of incidence $\cos \theta$.
- Momentum transferred per photon = $2 p_\nu \cos \theta$.

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Pressure = momentum flux

- **Pressure = outward (photons/sec/unit area) x (2 p_vcos θ)**
= (in+out)(photons/sec/unit area) x (p_vcos θ)

$$P_{\text{rad},v} = \frac{1}{c} \oint_{4\pi} I_v \cos^2 \theta \cdot d\Omega = \frac{2\pi}{c} \int_{-1}^1 I_v \mu^2 d\mu$$

- **The factor cos² θ allows for both foreshortening of the surface element's cross-section and transfer of normal component of momentum.**

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Radiative energy transport

- **Specific intensity of beam travelling at angle θ to radial direction in medium of density ρ, opacity κ_v and source function S_v:**

$$\cos \theta \frac{dI_v}{dr} = \rho \kappa_v (S_v - I_v). \text{ Multiply both sides by}$$

$\mu = \cos \theta$, and integrate:

$$\int_{-1}^1 \mu^2 \frac{dI_v}{dr} d\mu = \rho \kappa_v \int_{-1}^1 \mu d\mu (S_v - I_v) \quad S_v \text{ is isotropic}$$

$$\Rightarrow \frac{d}{dr} \int_{-1}^1 I_v \mu^2 d\mu = \rho \kappa_v S_v \int_{-1}^1 \mu d\mu - \rho \kappa_v \int_{-1}^1 I_v \mu d\mu$$

cf. Radiation pressure
cf. Flux

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Radiative-equilibrium temperature gradient

- We find that the opacity determines the temperature gradient:

$$c \frac{dP_{\text{rad},v}}{dr} = -\kappa_v \rho F_v \Rightarrow \frac{dP_{\text{rad}}}{dr} = -\frac{\kappa \rho}{c} F$$

where we define a flux - weighted mean opacity,

$$\kappa \equiv \frac{\int_0^\infty \kappa_v F_v dv}{\int_0^\infty F_v dv}. \text{ But } P_{\text{rad}} = \frac{u}{3} = \frac{1}{3} aT^4,$$

$$\text{so } \frac{dP_{\text{rad}}}{dr} = \frac{4}{3} aT^3 \frac{dT}{dr}. \text{ Also } F = \frac{L}{4\pi r^2},$$

$$\Rightarrow \frac{dT(r)}{dr} = -\frac{3\kappa\rho}{4acT^3} \frac{L(r)}{4\pi r^2} \quad \text{Another equation and a new variable, } T(r).$$

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Convective equilibrium

- Suppose temperature gradient is radiative.
- Is it stable to small local perturbations?
- Suppose a blob of mass δm at radius r has its temperature perturbed by a small amount:

$$\Delta T(r) = T_{\delta m}(r) - T(r).$$

- Pressure will change by

$$\Delta P(r) = P_{\delta m}(r) - P(r).$$

- but pressure balance is quickly restored by a change in volume, to give density difference from surroundings:

$$\Delta \rho(r) = \rho_{\delta m}(r) - \rho(r).$$

- Temperature excess with pressure equilibrium in ideal gas \rightarrow density deficiency \rightarrow buoyancy.

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Buoyant stability

- If $\Delta T > 0$, buoyant force > gravity, so blob rises to new position at $r+\Delta r$.

- Surroundings at new position have density

$$\rho(r) + \frac{d\rho}{dr} \Delta r$$

- while blob changes density to match local pressure:

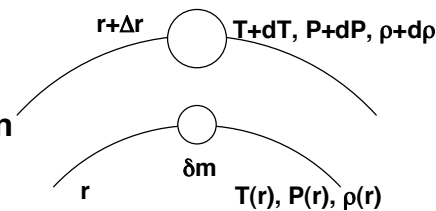
$$\rho_{\delta m}(r) + \left(\frac{d\rho}{dr}\right)_{\delta m} \Delta r$$

- Element is stable if it becomes denser than surroundings, i.e. if:

$$\left|\left(\frac{d\rho}{dr}\right)\right| > \left|\left(\frac{d\rho}{dr}\right)_{\delta m}\right|$$

- Ideal gas:

$$P \propto \rho T \Rightarrow \left|\left(\frac{dT}{dr}\right)\right| < \left|\left(\frac{dT}{dr}\right)_{\delta m}\right| \text{ for stability.}$$



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Adiabatic changes

- Rising blob is hotter than its surroundings as it rises so can only lose heat (& vice versa for falling blob).

- Hence change in temperature with r must be less than adiabatic (no heat loss) value :

$$\left|\left(\frac{dT}{dr}\right)_{\delta m}\right| < \left|\left(\frac{dT}{dr}\right)_{ad}\right|.$$

- Adiabatic gradient is given by:

$$PV^\gamma = \text{const} \Rightarrow P \propto \rho^\gamma \propto \left(\frac{P}{T}\right)^\gamma, \text{ where } \gamma = \frac{C_P}{C_V}$$

$$\Rightarrow P^{1-\gamma} T^\gamma = \text{const}$$

$$\Rightarrow (1-\gamma) \frac{dP}{dr} + \gamma \frac{P}{T} \frac{dT}{dr} = 0.$$

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Logarithmic T-P gradients

- For an adiabatic blob, we thus get:

$$\frac{1}{T} \left(\frac{dT}{dr} \right)_{\text{ad}} = \frac{\gamma - 1}{\gamma} \frac{1}{P} \left(\frac{dP}{dr} \right)$$

- or:

$$\left(\frac{d \log T}{d \log P} \right)_{\text{ad}} = \frac{\gamma - 1}{\gamma}.$$

- Use hydrostatic equilibrium to get a similar T-P relation for the radiative gradient:

$$\left(\frac{d \log T}{d \log P} \right)_{\text{rad}} = \frac{3\kappa L(r)P}{16\pi acT^4 GM}.$$

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Convective stability criterion

- Remembering that pressure is the same inside and outside blob at all times, we can write the stability criterion as:

$$\left| \left(\frac{d \log T}{d \log P} \right)_{\text{rad}} \right| < \left| \left(\frac{d \log T}{d \log P} \right)_{\text{ad}} \right|$$
$$\Rightarrow \frac{3\kappa L(r)P}{16\pi acT^4 GM} < \frac{\gamma - 1}{\gamma}.$$

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Convectively unstable regions

- **(1) Cores of massive stars:**

Radiation flux $L(r)/4\pi r^2$ can become very large while opacity $\kappa\rho$ remains small in the centres of main massive main - sequence stars.

- **(2) Outer envelopes of cool stars:**

Adiabatic exponent γ can approach unity in sub - surface ionization zones in cool stars.

Hence $(\gamma - 1)/\gamma$ can become small, and convection will set in at quite low values of $|(dT/dr)_{\text{rad}}|$.

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Energy transport

- **In formulating stellar structure problem, use a single expression for the temperature gradient:**

$$\frac{d \log T}{d \log P} = (1 - \xi) \left(\frac{d \log T}{d \log P} \right)_{\text{rad}} + \xi \left(\frac{d \log T}{d \log P} \right)_{\text{ad}} .$$

where ξ characterizes the convective efficiency :

$\xi = 0 \Rightarrow$ radiative equilibrium

$\xi = 1 \Rightarrow$ adiabatic convection

$0 < \xi < 1 \Rightarrow$ non - adiabatic convection : ξ must be determined from convection theory.

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2.5.1 Constitutive relations

- **Still need additional equations to describe ρ , ε , κ , ξ and $(d\log T/d\log P)$ in terms of:**
 - the state variables T and P, and
 - the composition of the stellar material (X,Y,Z or X_i)
- **The following *constitutive relations* close the system of ODEs:**

$$\rho = \rho(P, T, X_i) \quad (\text{equation of state})$$

$$\varepsilon = \varepsilon(\rho, T, X_i) \quad (\text{nuclear energy generation rate})$$

$$\kappa = \kappa(\rho, T, X_i) \quad (\text{opacity})$$

$$\xi = \xi(\rho, T, X_i) \quad (\text{convective efficiency})$$

$$\frac{d\log T}{d\log P}(\rho, T, \xi, \kappa, X_i) \quad (\text{energy transport})$$

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2.5.2 Equations of stellar structure

- **We have now determined the four basic (time-independent) equations of stellar structure.**
- **Use mass continuity to transform them to have enclosed mass as the independent variable.**

- **Mass continuity:**
$$\frac{dr}{dm} = \frac{1}{4\pi r^2 \rho}$$

- **Conservation of energy:**
$$\frac{dL}{dm} = \varepsilon$$

- **Hydrostatic equilibrium:**
$$\frac{dP}{dm} = -\frac{Gm}{4\pi r^4}$$

- **Energy transport:**
$$\frac{dT}{dm} = \frac{Gm}{4\pi r^4} \frac{T}{P} \frac{d\log T}{d\log P}$$

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2.5.3 Boundary conditions

- To solve a system of n ODEs, we need to specify n boundary conditions.
- In Lagrangian frame, boundaries are at the centre ($m=0$) and the surface ($m=M$).

At the centre:

$$r(m = 0) = 0$$

$$L(m = 0) = 0$$

At the surface :

$$T(m = M) = T_{\text{eff}} \left(= \frac{L}{4\pi R^2 \sigma} \right)^{1/4}$$

$$P_{\text{gas}}(m = M) = 0$$

2.5.4 Solution

- Solution of equations of stellar structure gives the run of P , T , m and L as functions of r throughout the domain $0 < r < R$.
- Solutions are characterized uniquely by
 - Total mass of star $M = m(R)$
 - Run of chemical composition through star.
 - Gravitational binding energy.
- Gives quantitative description of stellar interior.